

Constructing a Magnetic Tweezers to Monitor RNA Translocation at the Single-Molecule Level 2 3

Desiree Salas, Veronika Gocheva, and Marcelo Nöllmann 4

Abstract 5

Single-molecule methods have become an invaluable tool in the investigation of the mechanisms of nucleic-acid motors. Magnetic tweezers is a single-molecule manipulation technique that permits the real-time measurement of enzyme activities on single nucleic-acid molecules at high-resolution, high-throughput, and inherently constant force. Here, we describe several aspects of the implementation of magnetic tweezers, with special emphasis on the construction of a simple magnetic trap and, in particular, on the detailed description of image analysis methods to measure the extension changes in nucleic-acid molecules induced by protein activity. Finally, we carefully describe the steps involved in performing a full magnetic tweezers experiment. 6 7 8 9 10 11 12 13

Key words Motors, Nucleic-acid enzymes, RNA, Magnetic tweezers, Single molecule, Mechanochemistry 14 15

1 Introduction 16

Force has an essential role in a myriad of biological mechanisms. For example mechanical forces are generated in the cell during several processes like replication, transcription, translation, chromosomal segregation, and many others. Several techniques have been developed in the last 20 years with the capability to measure the forces and displacements generated by single molecules or to directly apply external forces to these processes in order to reveal their underlying molecular mechanisms. Single-molecule nano-manipulation techniques have made it possible to study protein activity under an applied mechanical stress. In fact, large structural rearrangements are often rate limiting, and applying assisting or opposing force can modulate these conformational changes and hence the catalytic cycle of molecular motors. The three commonly used single-molecule force-based manipulation methods are atomic force microscopy, optical tweezers, and magnetic tweezers. Here, we will describe the implementation of magnetic tweezers to directly measure the translocation activity of RNA motors. 17 18 19 20 21 22 23 24 25 26 27 28 29 30 31 32

33 In a typical magnetic tweezers setup, a single nucleic-acid (NA)
34 molecule is specifically tethered to a surface on one end and to a
35 paramagnetic bead on the other (Fig. 1a). A pair of external perman-
36 ent magnets generates a magnetic field that induces a magnetic
37 moment on the bead that aligns with the lines of field (Fig. 1a). This
38 induction leads to an external force that is proportional to the gradi-
39 ent of the scalar product between the induced magnetic moment and
40 the magnetic field. This force is perpendicular to the field lines
41 and points in the direction of the magnets (Fig. 1a). As the bead is
42 tethered to the surface by single molecules, the magnetic force traps
43 the bead in the axial direction [1, 2]. The force can be controlled by
44 vertically translating the magnets in the axial direction and decreases
45 exponentially with a typical decay constant of ~ 1 mm [2, 3]. Typical
46 forces used in magnetic tweezers range from a few fN to ~ 10 – 20 pN
47 [2], although larger forces can be attained (~ 100 pN). Several meth-
48 ods have been developed to estimate the trapping force by analysis
49 of the Brownian motion of trapped beads [1, 3, 4].

50 In magnetic tweezers, the distance between the bead and the
51 surface can be used as a direct, instantaneous measure of the NA
52 tether length and employed to determine changes in NA extension
53 due to protein activity (Fig. 1b). Importantly, the changes in NA
54 tether extension (\sim few μ m) are considerably smaller than the decay
55 constant of the magnetic field (~ 1 mm), and therefore, the force
56 on the paramagnetic bead can be considered constant and only
57 defined by the distance between the magnets and the coverslip.
58 The relation between extension of the tether and applied force
59 depends on the mechanical properties of the NA molecule [5].
60 These mechanical properties vary considerably between single-
61 stranded [6] and double-stranded DNA and RNA [7–10], and this
62 difference was often used to monitor the conversion between sin-
63 gle- and double-stranded NA molecules [11]. Importantly, mag-
64 netic tweezers allow for the application of torque by simply rotating
65 the magnets [8] and the measurement of torque introduced by the
66 activity of enzymes [12–15]. Critically, the ability of magnetic
67 tweezers to apply torque on single double-stranded, topologically
68 closed DNA molecules permits the introduction and control of the
69 degree of supercoiling, making it the method of choice to investi-
70 gate, for instance, the mechanisms of supercoiling regulation by
71 DNA topoisomerases [16–19].

Fig. 1 (continued) extension of the NA tether is reduced (*right panel*). **(b)** The activity of RNA motors can lead to the change in the extension of the tether either by looping of the RNA substrate or by conversion from double- to single-stranded RNA (helicase activity). **(c)** The basic components of the illumination system and magnetic trap are shown: light from a fiber is first filtered by color to make it monochromatic, then collimated and guided to the sample by lenses L_1 , L_2 , and L_3 (see text). Magnets are positioned so that their axis is colinear with the optical axis of the objective and the illumination system. The position of the objective in the optical axis is controlled by a piezoelectric stage

RNA Translocation Monitored by Magnetic Tweezers

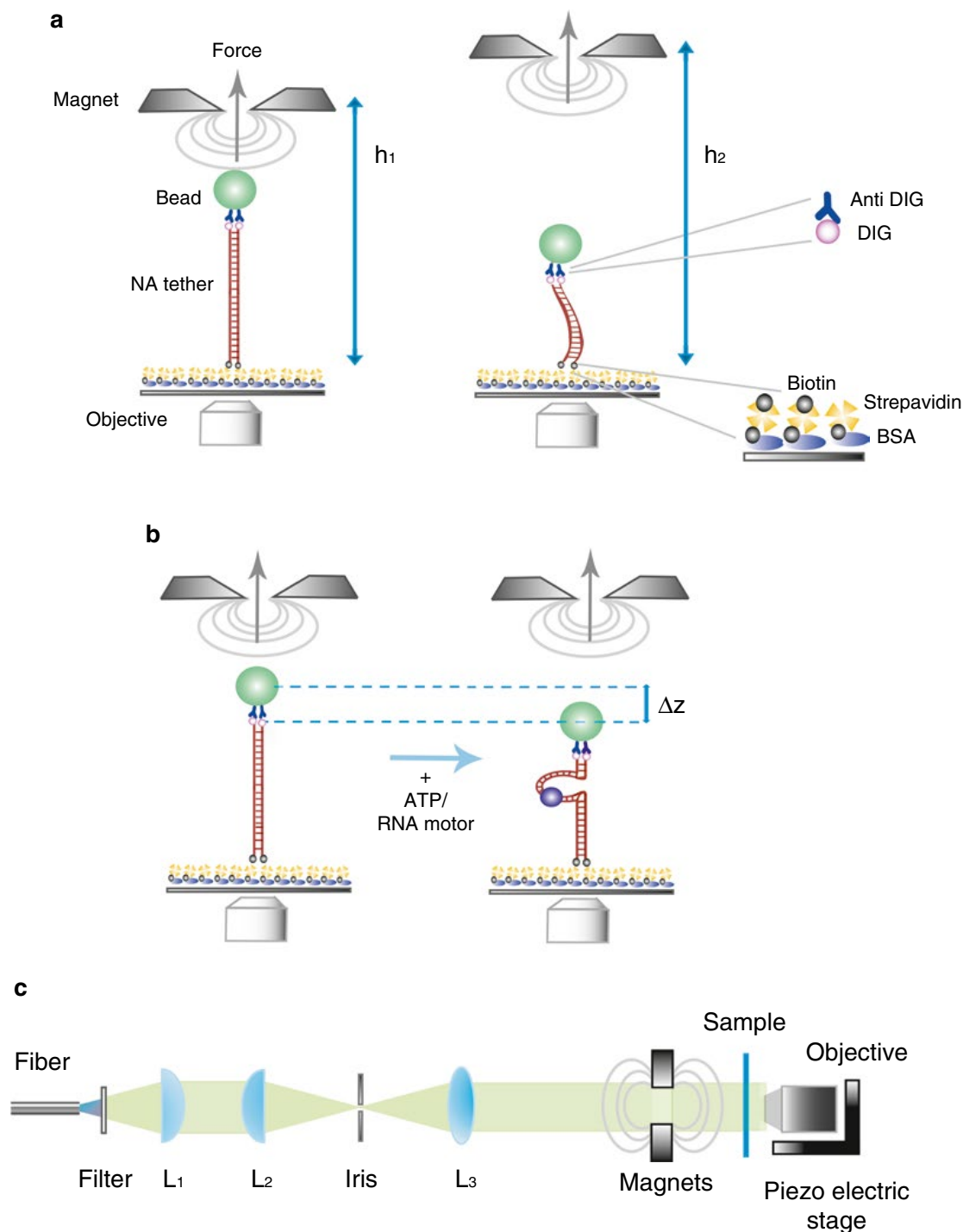


Fig. 1 Experimental design. **(a)** Schematic representation of a magnetic tweezers setup. A paramagnetic bead (green sphere) is attached to a single nucleic-acid (NA) molecule (red vertical lines) by DIG-anti-DIG interactions (pink sphere and blue tripod, respectively). On the other end, the NA tether is attached to the surface of the coverslip a flow cell by BSA-biotin/streptavidin/biotin linkages (yellow trapezoids, gray sphere, respectively). The presence of a pair of small magnets produces a magnetic field (gray lines) that produce a force on the bead proportional to the gradient of the field. The magnetic force exponentially decays with the distance between the magnets and the observation chamber (h). At higher distances, the force decreases and the

72 Magnetic tweezers are based on video microscopy and thus
73 have an important advantage over other manipulation methods as
74 multiple, single NA/bead attachments can be visualized at the
75 same time. This advantage has been recently used to simultane-
76 ously monitor the trajectories of tens to hundreds of single mole-
77 cules, considerably increasing the high throughputness of the
78 method [20–22] and uniquely allowing for the detection of rare
79 events at the single-molecule level [23].

80 The construction of a magnetic tweezers setup to study RNA
81 helicases involves many hardware and software developments.
82 Several methodological publications describe in detail the proce-
83 dures used for the fabrication and use of microfluidics chambers
84 for magnetic tweezers [24, 25], for the design and construction of
85 DNA [25, 26] and RNA handles [10, 27], and for the character-
86 ization of forces [1, 3, 4]. In this chapter, we will concentrate on
87 the construction of a simple magnetic trap and on the detailed
88 description of the methodology used for image analysis.

89 Most algorithms used to determine the axial position of the
90 magnetic bead rely on the interference ring pattern created by out-
91 of-focus beads [28]. This interference pattern can be used to track
92 beads in 3D with high precision. Several algorithms have been
93 described in the literature that directly compare images of beads
94 with simulated or real images of a library to obtain a vertical reso-
95 lution of ~ 10 nm [2, 29–31]. Here, we will describe a considerably
96 simpler method that does not rely on a comparison to a library of
97 images, but rather on interpolating a signature of the interference
98 ring that can be directly linked to the position of the bead in the
99 vertical axis.

100 The methods section will begin by the construction of a simple
101 magnetic trap, continue with a description of the image analysis
102 methods used to extract the three-dimensional trajectory of mag-
103 netic beads (which translates into a measure of the extension of the
104 NA tether), and end with a discussion of the steps involved in per-
105 forming a typical magnetic tweezers experiment.

106 2 Materials

107 2.1 Illumination 108 System

- 109 1. Fiber white light source (Thorlabs, OSL1-EC or OSL2).
110 Lenses: L1 and L2 as aspheric lenses ($f=17$ mm,
111 MAP052525-A1), L3 is a spherical lens ($f=10$ cm).
- 112 2. Translation stage (Thorlabs, PT1) mounted on a 1.5" mounting
113 post system (Thorlabs, P350/M, PB4/M, C1525, and PF175).
- 114 3. Cage components (30 mm diameter): four rods length 20 cm
115 (Thorlabs, ER10) and material to mount cage to translation
stage (Thorlabs, UPH6, TR3).
4. Bandpass filter (550–600 nm).

2.2 Microscope Components

1. An active optical table with vibration control (e.g., VH3660W-OPT table from Newport). 116
117
2. Commercial microscope with 1.4 numerical aperture, 100× oil objective, and at least one C-mount port. 118
119
3. CMOS (complementary metal oxide semiconductor) or CCD (charge-coupled device) camera (e.g., Neo and DU-885 cameras from Andor are good options for coupling with fluorescence applications). The camera has to be able to perform fast exposures (faster than 1 ms). 120
121
122
123
124
4. Closed-loop, high-resolution piezoelectric stage (MCL NanoF100). 125
126
5. Closed-loop motorized XY translation stage with rotary or linear encoders (ASI S31121010FT). The key feature in this component is the long-term stability, low lateral drift, and low degree of coupling between axes. 127
128
129
130
6. Digital to analog converter (NI USB-6211) (*see Note 1*). 131
7. Micrometer ruler slide (e.g., R1L3S2P from Thorlabs). 132

2.3 Magnet Mount

1. High-resolution linear translation stage (PI, M404.4PD) (*see Note 2*). 133
134
2. Fast rotation stage (typically >4 Hz) with high repeatability (PI, M-660.55) to permit the measurement of rapid, enzyme-induced changes in supercoiling levels and allow for accuracy in the angular positioning of the magnet axis. 135
136
137
138
3. 1.5" mounting post system (Thorlabs, P350/M, PB4/M, CI525, and PF175) mounted on crossed, manual linear micrometric stages (Thorlabs, PT1) to align magnets on the optical axis. 139
140
141
142

2.4 Reagents and Small Equipment

1. Dynabeads MyOne carboxylic acid beads, 1 μm (Life Technologies) (*see Note 3*). 143
144
2. Phosphate buffer saline (PBS): 137 mM NaCl, 2.7 mM KCl, 10 mM Na₂HPO₄, and 1.8 mM KH₂PO₄. 145
146
3. Dynabeads MyOne streptavidin C1 beads (Life Technologies). 147
4. Streptavidin-coated magnetic beads 1 μm (Invitrogen). 148
5. Anti-digoxygenin (DIG) antibody. 149
6. Bovine serum albumin (BSA). 150
7. Biotinylated BSA solution: biotinylated BSA at 5 mg/ml in PBS. The solution is prepared from lyophilized biotinylated BSA by dilution into PBS (5 mg/ml), vigorous vortexing, and centrifugation. The solution is divided into 20 μl aliquots and stored at -20 °C. 151
152
153
154
155
8. 2 mg/ml streptavidin in PBS. Store at -20 °C as 30 μl aliquots. 156
157

- 158 9. 60 pmol/ μ l NA tether stock solution [10, 25–27] (see also
159 chapter by Wei Cheng in this book).
160 10. Binding buffer: 50 mM Tris–HCl pH 7.5, 1 M NaCl, 1 mg/
161 ml BSA.
162 11. Rotary incubator.

163 2.5 Microfluidics

- 164 1. Home-made microfluidics chamber [24, 32]. Coverslips are
165 typically washed with ethanol and dried before assembly of the
166 microfluidics chamber. A detailed protocol for the preparation
167 of microfluidic chamber is provided in the chapter by Wei
Cheng in the present volume.

168 3 Methods

169 All steps below are performed at room temperature unless specified
170 otherwise.

171 3.1 Construction 172 of a Magnetic Trap

- 173 1. Assemble the illumination system as described in Fig. 1c. The
174 first lens (L1) collimates the beam from the fiber light source,
175 the second (L2) focuses it on the iris diaphragm, while the third
176 lens (L3) is away from the iris by its focal length and collimates
the beam into the flow chamber (Fig. 1c). Place the cage system
vertically and align the fiber light source (*see Note 4*).
177 2. Magnets are centered on the optical axis of the objective
178 (*see Note 5*).
179 3. Measure the pixel size of the camera by using a micrometer
180 ruler slide. Typical pixel sizes are 80–100 nm.

181 3.2 Preparing 182 Magnetic Beads 183 and Generating 184 NA-Bead 185 Attachments

- 186 1. Assemble the microfluidics chamber and flush 1 ml of PBS
187 through the channel.
188 2. Replace PBS in the channel with biotinylated BSA solution (s).
189 Incubate for 2 h (*see Note 6*).
3. Flush out the BSA-biotin solution by passing 1 ml of PBS.
4. Prepare a fresh 0.2 mg/ml solution of streptavidin (*see Note*
7) and flush it into the channel. Incubate for 10 min.
5. Flush out streptavidin by passing 1–2 ml of PBS through the
channel.
6. Transfer 0.5 μ l of NA tether stock solution (60 pmol/ μ l) into a
microtube and add 1 μ l of a 1:10 dilution of anti-DIG beads
(*see Note 3*). Add 1 μ l of a dilution of streptavidin magnetic
beads to the mix to be used as reference beads for drift correc-
tion. At the final concentration, streptavidin beads should be at
a 1:10 molar ratio with respect to anti-DIG beads. Avoid pipet-
ting this mix solution back and forth to limit NA shearing.

- | | | |
|---|--|--|
| 7. | Add 3.5 μl of binding buffer and mix quickly by vigorous hand shaking. Avoid vortexing or centrifugation at all price, since these will shear the NA tether and precipitate beads. | 197
198
199 |
| 8. | Incubate solution in a rotary incubator for 10 min at room temperature. Then, add 45 μl of binding buffer. | 200
201 |
| 9. | Flush the solution into the channel of the microfluidic chamber while continuously imaging the coverslip surface on the CCD. Stop flow when beads appear on the field of view (FOV). After a short period, beads should start to fall on the coverslip surface. | 202
203
204
205
206 |
| 10. | After ~ 10 – 20 beads have set per FOV, turn the flow back on. The flow rate should be such that beads move on the surface at $\sim 5 \mu\text{m/s}$ (see Note 8). Magnets should be far from the surface (> 5 – 10 cm) during incubation; otherwise, beads will aggregate on the top surface of the channel. | 207
208
209
210
211 |
| 11. | After 10–20 min, approach the magnets to ~ 2 – 3 cm from the surface and explore different FOVs to detect proper bead attachments. A further description of this process can be found elsewhere [24]. | 212
213
214
215 |
| 12. | Wash unattached beads until no more beads are seen to flow in the chamber. Only beads attached to the coverslip should be present. Nonspecifically attached beads can often be removed by manually passing a pair of fridge magnets close to the top surface of the microfluidics chamber. | 216
217
218
219
220 |
|
 | | |
| 3.3 Obtaining the Height of the Bead from Image Analysis | 1. After illuminating the sample, image the interference pattern created by out-of-focus beads on a CCD camera. A typical FOV with several beads is shown in Fig. 2a. Beads are automatically segmented by following the procedure described in steps 2–6. | 221
222
223
224
225 |
| | 2. Subtract the minimum intensity of the image from all pixel intensity values and scale the resulting image so that its maximum intensity is unity. | 226
227
228 |
| | 3. Threshold the image to obtain a binary image. The threshold can be obtained by using the Otsu method, an algorithm that calculates the optimum threshold so that the combined intra-class variance between the two main intensity levels in the image is minimal. Perform a flood-fill operation on the background pixels to eliminate holes in the binary image. | 229
230
231
232
233
234 |
| | 4. Remove objects in the binary image with less than eight pixels (this parameter can be adjusted to improve results). | 235
236 |
| | 5. Classify objects in the binary image by finding the connected components. | 237
238 |
| | 6. Determine the centroids (x_i, y_i) and areas A_i of each object i . Keep only objects with large areas (usually > 500 pixels). | 239
240 |

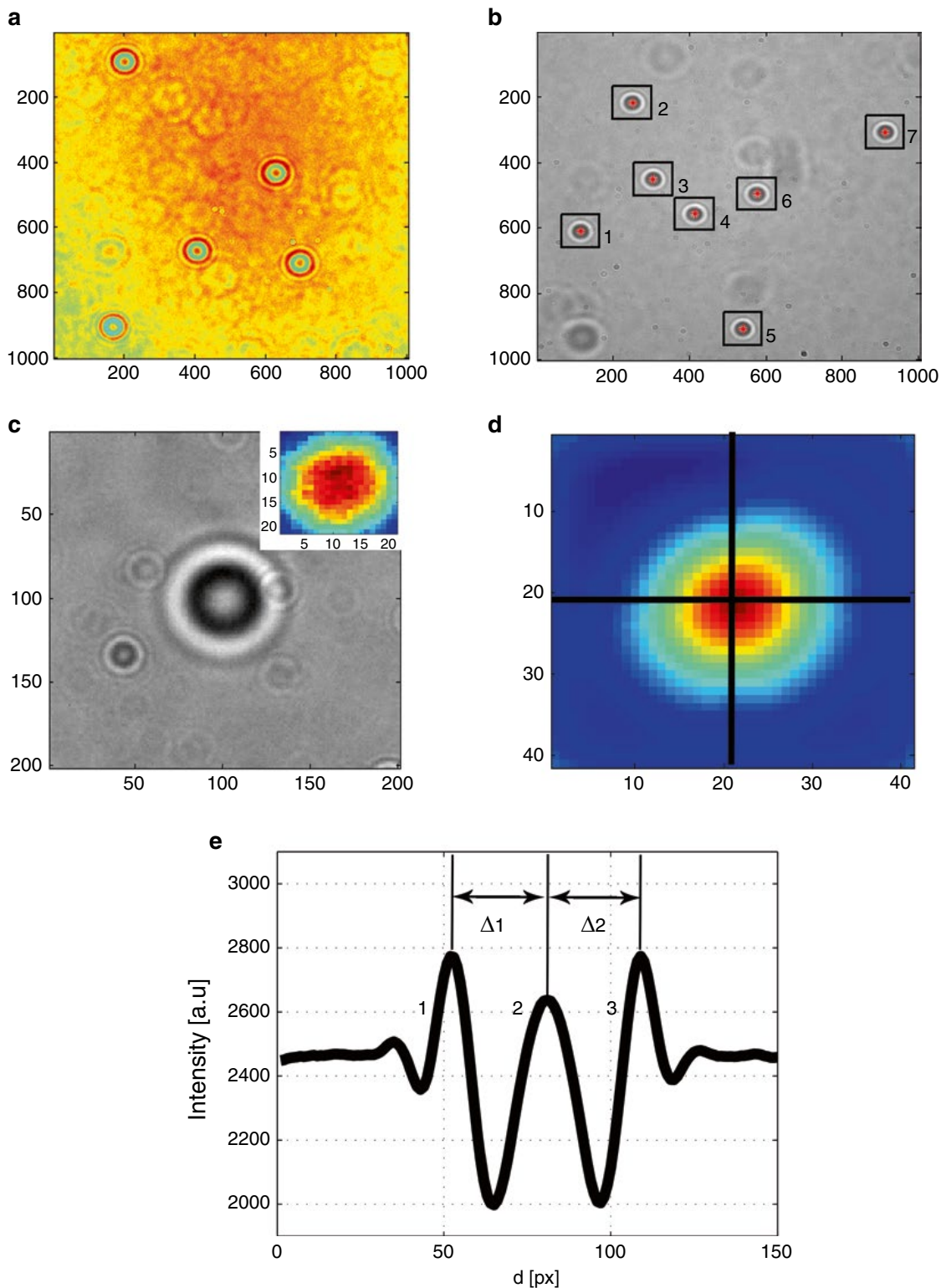


Fig. 2 Measurements of the bead axial position. (a) Interference patterns displayed by out-of-focus paramagnetic beads ($1 \mu\text{m}$) tethered to the surface by single NA attachments. Axes are in pixel units. (b) Bead segmentation is performed to automatically detect beads in the FOV. Boxes delimiting the ROIs used to extract the ring patterns of the beads (200×200 pixels) are shown as *black boxes*, while the position of the centers of the ROIs

- The results of a typical segmentation are shown in Fig. 2b. Boxes around each bead represent the region of interest (ROI) occupied by that bead (usually 200×200 px, but this number depends on the total axial range to be calibrated and the density of beads in the FOV).
7. A typical ROI is shown in Fig. 2c (this image will be hereafter called I_1). The out-of-focus bead displays clear symmetric rings. Importantly, the distance between the rings and the center of the bead depends on the distance of the bead to the focal plane. This property is used to determine the axial position of beads. The first step in this procedure involves determining the centroid of the bead (center of symmetry). Our preferred method involves image cross-correlation and is described in steps 8–11.
 8. Crop the image of the bead further by using a smaller ROI (usually 20×20 pixels). An example of the result of this operation is shown in the inset of Fig. 2c and will be hereafter called I_{ROI} .
 9. Flip I_{ROI} in the horizontal and vertical directions to obtain $I_{ROI-flipped}$.
 10. Determine the normalized cross-correlation between I_{ROI} and $I_{ROI-flipped}$ (I_{CC}) (see Fig. 2d for the cross-correlation function from the inset of Fig. 2c). Note that I_{CC} will have two times the vertical and horizontal dimensions of I_{ROI} .
 11. Find the maximum of I_{CC} by determining the pixel with maximum intensity. First, determine the intensity profiles in the x and y directions at the center of the maximum. To interpolate the maximum of the I_{CC} function with sub-pixel resolution, perform a quadratic polynomial fit separately on each intensity profile and determine analytically the coordinates of the I_{CC} maximum from the coefficients of the polynomial series. Divide these coordinates by two to obtain the coordinates of the centroid of I_{ROI} with sub-pixel resolution. Convert, these coordinates into the coordinates of the image of the bead I_1 (Fig. 2c) by taking into account the different sizes of I_1 and I_{ROI} .
 12. Determine the distance of the first concentric ring to the centroid of the bead as described in steps 13–15.

←
Fig. 2 (continued) are shown in *red*. Axes are in pixel units. **(c)** A typical ROI in which the interference rings can be clearly observed (bead 1 in panel **b**). Inset shows the cropped image (21×21 pixels) used for image cross-correlation (see text). Axes are in pixel units. **(d)** Image cross-correlation the inset in panel **c**. Crossed hairline indicates the center of the ROI. Coordinates of the bead centroid are calculated from the interpolated maximum of the cross-correlation image (see text). **(e)** Average intensity profile (ring pattern) of the intensity signal across the centroid of the bead (from image in panel **c**). Distances between peaks 1–2 and 2–3 define the values of Δ_1 and Δ_2

- 278
279
280
281
282
283
284
285
286
287
288
289
290
291
292
13. Using the coordinates of the centroid of the bead in image I_1 (**step 11**), plot the intensity profiles in the x and y directions across the centroid. The two intensity profiles are averaged to obtain the approximated ring profile of the bead (Fig. 2e, black thick line). Three maxima are clearly visible that correspond to the central bright spot of the bead and the first bright ring.
 14. Determine the positions of the peaks with sub-pixel resolution by looking for downward zero crossings that exceed a certain threshold (usually set to zero).
 15. Determine the distances Δ_1 and Δ_2 between the peaks corresponding to the centroid of the bead (peak 2 in Fig. 2e) and the first concentric ring (peaks 1 and 3). Average Δ_1 and Δ_2 to obtain Δ (pixel units). Δ is a parameter that shows a one-to-one correspondence to the axial distance between the bead and the focal plane (*see Note 9*).

293
294
295
296
297
298
299
300

Other algorithms exist that obtain a parameter from the image that is strictly related to the vertical position of the bead. These algorithms use the Hilbert transform to calculate the phase shift between ring patterns [2] or calculate the entropy of the image without even needing to obtain a precise estimation of the centroid of the bead. In our hands, the algorithm described here is the fastest and displays excellent vertical resolution (~ 1.4 nm in a 1 s window; see Fig. 3d).

301 **3.4 Determining**
302 **and Applying**
303 **a Calibration**

- 304
305
306
307
308
309
310
311
312
313
314
315
316
317
318
319
320
321
1. To determine the relation between Δ and the distance of the bead to the objective focal plane (calibration function), take images of the bead at different focal positions of the objective using the piezoelectric stage holding the objective (usually steps of 50 nm are taken) (*see Note 10*). The ring profile dramatically changes with the position of the objective (defined as z) (Fig. 3a). In the example shown, six positions were acquired with a step size of 500 nm.
 2. Plot the value of Δ for each profile i (Δ_i) against z_i (Fig. 3b). Use a second-order polynomial ($z(\Delta) = A_0 + A_1 \Delta + A_2 \Delta^2$) to interpolate Δ for any value of z in the scanned interval (solid line in Fig. 3b).
 3. Test the calibration by first acquiring a dataset in which the whole z range is scanned. In a test example, we obtained ten images for each z position, with a step size of 500 nm. The resulting images were analyzed as described in Subheading 3.3 and Δ was obtained for each frame. Δ was then converted into height by using the polynomial interpolation $z(\Delta)$ (**step 2**). The obtained z positions were plotted against frame number to verify that the distance between steps corresponds to the step size used in the acquisition (Fig. 3c).

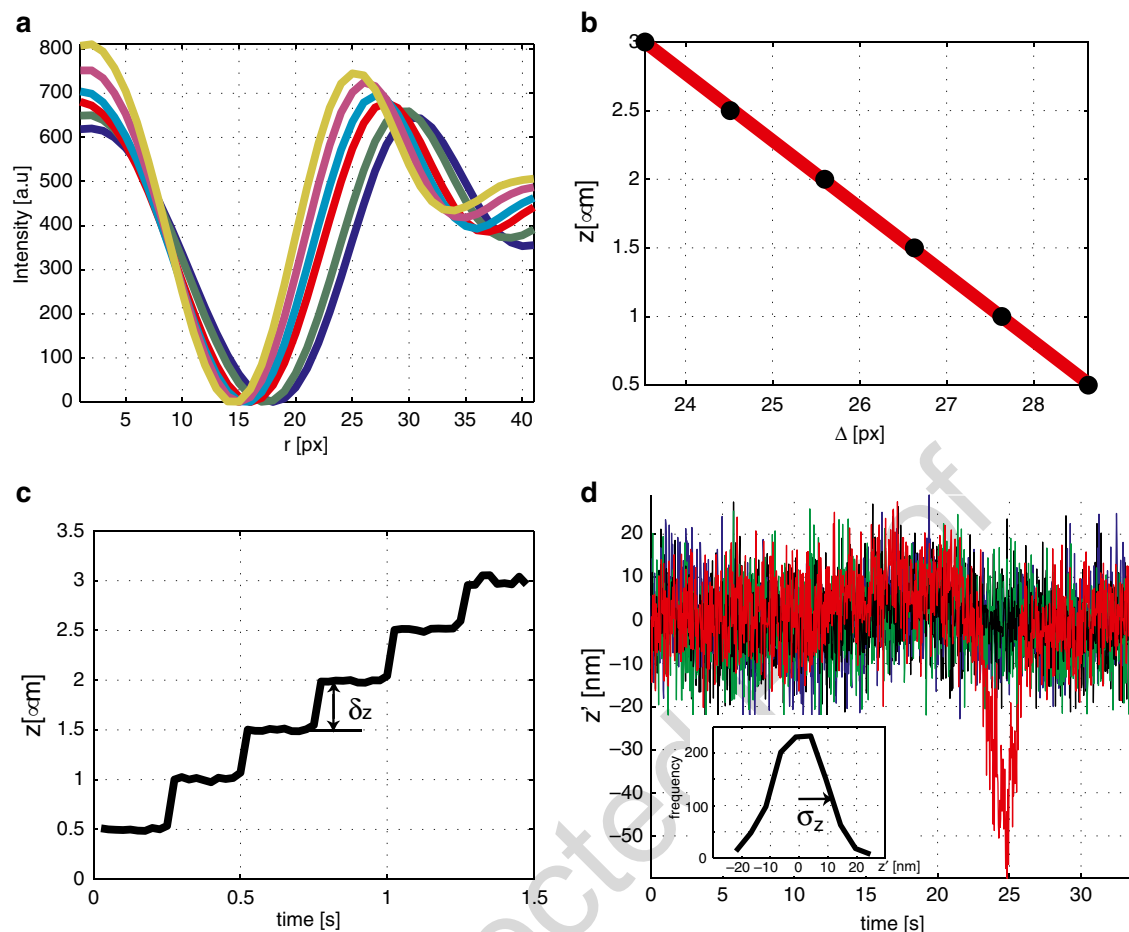


Fig. 3 Application of the bead calibration. **(a)** Variation of the ring pattern with position of the focal plane. The change in phase of the rings is approximately proportional to the distance of beads to the focal plane. **(b)** Plot of the focal position (z) as a function of Δ (black open circles). Solid line represents a second-order polynomial fit to the experimental data. **(c)** The calibration is tested by running a mock experiment in which the focal position of the objective is shifted in steps of $\delta z = 500$ nm (10 images per position). **(d)** Single-molecule time traces of the vertical position of four tethered beads in the same FOV (z'). Most tethers do not show activity (black, blue, and green) and their vertical position fluctuates due to Brownian motion. In contrast, the red trace shows a typical single-molecule ssRNA looping event. Inset shows the frequency of z' positions for a single-tethered bead displaying Brownian motion. σ_z represents the standard deviation of the distribution (10 nm at 40 Hz)

For an FOV with multiple beads, calibration parameters are 322
 obtained for each bead and used to analyze the distance changes of 323
 each bead during an experiment. To perform this process in teth- 324
 ered beads, high forces (~ 10 pN) are used to limit axial movement 325
 during the acquisition of the calibration file. 326

3.5 Performing an Experiment and Analyzing Data

1. After assembling the microfluidics chamber and obtaining 327
 tethered beads (procedure extensively described elsewhere 328
 [24]), characterize the length of each attached bead in the 329
 FOV, using the buffer in which experiments will be performed. 330

331 The most common procedure involves measuring the tether
 332 extension as a function of force and performing a fit to a worm-
 333 like chain model. This procedure provides the persistence length
 334 of the NA tether and its full extension. Use only beads displaying
 335 correct full extension and persistence length for experiments.

- 336 2. After tether characterization, introduce a solution containing
 337 the protein under study into the chamber (*see Note 11*).
- 338 3. Any protein activity that changes the length of the NA tether
 339 (like a looping enzyme that shortens the extension of the tether,
 340 or a helicase converting the NA construct from double-stranded
 341 to single-stranded form) will generate a change in the vertical
 342 position of the bead. The change in axial position of each bead
 343 in the FOV is measured by determining Δ for each bead and
 344 each image and by using the bead-specific calibration function
 345 to obtain a z position. This calculation can be performed in off-
 346 line mode (by registering a movie of the FOV over long time
 347 periods and performing post-analysis, *see Note 12*) or in real
 348 time (by image analysis performed on-the-fly, *see Note 13*).
- 349 4. Correct z -axis values for the displacement of the objective. The
 350 displacement of the objective does not correspond to that of
 351 the focal plane due to refraction of light at the water-glass
 352 interface. The real axial position of the bead (z') is obtained by
 353 $z' = z n_{\text{water}} / n_{\text{glass}} \sim 0.86 z$, where n_{water} and n_{glass} are the indexes
 354 of refraction of water and glass and z is the axial position calcu-
 355 lated by moving the piezoelectric stage.
- 356 5. Correct the extension over time for each bead, $z'_i(t)$, to account
 357 for axial drift (*see Note 14*). This procedure is performed by fol-
 358 lowing in parallel the z position of a bead that is stuck to the
 359 surface ($z'_{\text{stuck}}(t)$; *see Note 15*) and by subtracting the average
 360 axial position of the stuck bead from the positions of all beads (*see*
 361 *Note 16*). In other words, drift corrected extensions are obtained
 362 by $z'_{i, \text{corrected}}(t) = z'_i(t) - \langle z'_{\text{stuck}}(t) \rangle$. An example of the tracking
 363 of several beads after drift subtraction is shown in Fig. 3d.
- 364 6. Scan the traces of each bead in the FOV manually or automati-
 365 cally to detect single-molecule events (*see Note 17*). A typical
 366 looping event is shown in Fig. 3d. These events can be used to
 367 obtain information on the processivity, the velocity, and the
 368 pause length behavior of the enzymatic reaction at the single-
 369 molecule level.

370 4 Notes

371 1. A DAQ is essential for performing the z -axis calibration by
 372 triggering the camera exposure at the same time as the piezo-
 373 electric stage is moved. If communication with the piezoelec-
 374 tric stage is through an analog signal, the card should have

- enough bits (usually 16 bits is recommended) to encode small enough changes in the position of the stage. 375
376
2. This translation stage should be fast (>10 mm/s) and precise (usually $\sim 1\text{--}2$ μm) to allow for rapid and small changes in force, show a low-pitch yaw angle (~ 50 μrad) to avoid lateral displacement during a long vertical translation, display high repeatability (typically $\sim 1\text{--}2$ μm) to ensure repeatable forces are attained during an experiment, and accept high loads (>100 N) in vertical configuration to allow for the displacement of magnets and the rotation stage. 377
378
379
380
381
382
383
384
 3. MyOne beads are functionalized with anti-DIG antibodies by following the standard protocol of the vendor. 300 mg of antibody are used per derivatization procedure. Derivatized beads can be kept at 4 $^{\circ}\text{C}$ for years. Beads will precipitate to the bottom of the tube and form large aggregates. Before use, the stock is vortexed and 2 μl of bead slurry are taken to make a 1:10 dilution of the stock in a 0.5 ml microtube using PBS buffer (final volume 20 μl). This dilution is sonicated by placing the microtube in a glass container filled with water and by approaching the microtube to a tip sonicator (for a total time of $\sim 5'$). This same 1:10 solution can be used two to three times (sonication tends to shear beads and so it is not desirable to sonicate the same bead solution many times). 385
386
387
388
389
390
391
392
393
394
395
396
397
 4. It is important that the lamp is placed on a shelf and not in contact with the optical table to avoid the transmission of vibrations to the microscope. 398
399
400
 5. Magnet alignment is achieved by positioning the magnets close to the objective lens and imaging the resulting transmitted light pattern on the camera. The magnets act as an iris that blocks the light transmitted to the objective. Their alignment on the optical axis is performed by maximizing the symmetry and homogeneity of the transmitted light pattern on the camera by displacing the micrometer manual translation stages used to mount the magnets. Then, the magnets are displaced vertically, and the process is repeated to ensure that the axis of translation of the vertical translation stage holding the magnet is colinear to the optical axis. Once tethers are obtained, a finer alignment is made by rotating the magnets while simultaneously tracking the centroid of tethered beads. Misaligned magnets lead to non-vertical forces and to elliptical orbits of the tracked bead during magnet rotation [33]. 401
402
403
404
405
406
407
408
409
410
411
412
413
414
415
 6. For each channel, a new tube of biotinylated BSA is thawed and used only once. By inserting air before and after the arrival of BSA-biotin into the channel, the use of this relatively expensive reagent can be minimized, and the arrival of the BSA-biotin solution can be directly observed. Tubes are clamped 416
417
418
419
420

- 421 during the incubation to avoid water evaporation and drying
422 of the channel.
- 423 7. For each experiment, an aliquot of streptavidin stock is thawed
424 and a 1:10 dilution is made with PBS to a final volume of
425 50 μ l. The dilution can be kept at 4 °C for no more than
426 1 week (do not refreeze).
- 427 8. If the flow rate is too high, beads will not remain in contact to
428 the surface and attached beads may nonspecifically bind to the
429 surface. If flow is too slow, beads risk on bouncing on each
430 other to form large aggregates.
- 431 9. The determination of bead height by using Δ is robust, as (1)
432 Δ is independent on intensity, and (2) the fact that the relative
433 distance of the first peaks to the center are used (and not the
434 real position of the center) means that errors in the centroid
435 determination show a lower degree of coupling into the height
436 determination than with other methods that rely on the abso-
437 lute ring pattern.
- 438 10. Calibration curves are obtained for each tethered bead. To
439 ensure that the beads move as little as possible in the axis due
440 to Brownian motion during the acquisition, the force exerted
441 by the magnet is set to its maximum (usually \sim 10 pN). In addi-
442 tion, the exposure time of the camera is reduced as much as
443 possible to limit image blurring (typical 2–5 ms).
- 444 11. The enzyme is prepared at a low enough concentration to
445 ensure single-molecule conditions (typical concentrations are
446 1–10 nM, but depend on the affinity of NA association con-
447 stant of the enzyme and on its oligomerization constant).
448 Typically, the buffer used is the same as that providing strong
449 “in bulk” activity, but usually contains BSA to avoid nonspe-
450 cific interactions. Often, enzymes show the tendency of inter-
451 acting nonspecifically with the tubes and glass in the chamber
452 (even if these are passivated by BSA). This effect can be more
453 deleterious at low enzyme concentration as, in this case, most
454 of the protein may end up in the tubes and not in the NA sub-
455 strate. In such cases, a low concentration of a detergent (i.e.,
456 Tween-20) or a higher concentration of BSA (0.1 mg/ml
457 final) can be used. Protein is usually kept on ice before injec-
458 tion into the chamber to maintain its enzymatic activity.
- 459 12. For off-line acquisition and analysis, movies are acquired in
460 frame transfer mode, with a 25 ms exposure time for as long as
461 30 min. Data are stored in spooling mode on a solid-state drive
462 (SSD) to improve writing and reading speeds (the latter is very
463 important for image analysis of long movies). The SSD disk is
464 emptied before each experiment as writing large files in spool-
465 ing mode in fragmented drives often leads to incomplete mov-
466 ies or movies with empty frames. Movies are saved as

- uncompressed TIFF files to avoid data loss during compression. Image analysis in MATLAB is performed on the movies to track beads in 3D and analyze trajectories. 467
468
469
13. Real-time acquisition is paramount when the user needs to react to the activity of the enzyme. Typical example is the study of topoisomerases, in which supercoiling needs to be introduced to generate the substrate on which the enzyme reacts (see for example [19]). After a processive reaction cycle of the enzyme, a new substrate needs to be generated by the user, making a real-time analysis tool almost unavoidable. A second case in which real time is essential arrives when the activity of the enzyme destroys the substrate, as is the case for DNA translocases that can irreversibly pull DNA to the surface unless sufficient force is exerted on the bead. As the initiation rate of the translocation activity is strongly force dependent, a real-time feedback on the current activity of the motor protein is required (see for instance [34]). We implemented real-time acquisition and analysis in LabVIEW with analysis routines implemented as external C libraries. 470
471
472
473
474
475
476
477
478
479
480
481
482
483
484
485
 14. Typical magnetic tweezers experiments rely on following the extension of tethered beads for extended periods of time (10 min to hours). During these long periods, thermal expansion or contraction of the different components of the microscope due to temperature changes in the environment produces slow changes in the distance between the objective and the chamber (drift). This drift can be minimized by properly choosing the components used in the microscope, assuring a constant temperature in the room and of all components used (such as the buffer injected into the microfluidics chamber), but it is ultimately impossible to eliminate it completely. 486
487
488
489
490
491
492
493
494
495
496
 15. In FOV with many beads, it is common to find beads attached nonspecifically to the surface that can serve as stuck beads for drift subtraction. Alternatively, the bead-NA mix can be supplemented with a low concentration of streptavidin beads that will specifically fix to the surface. A third possibility is to fuse polystyrene beads to the surface by depositing them and heating them on the coverslip before assembly of the microfluidics chamber. 497
498
499
500
501
502
503
 16. The drift of the chamber during acquisition can be obtained by following the axial movement of a stuck bead $z'_{\text{stuck}}(t)$ and filtering the signal by using a mean-filter function that averages the signal on a defined time window (typically ~ 1 s) to obtain $\langle z'_{\text{stuck}}(t) \rangle$. The drift in all beads can be corrected by subtracting $\langle z'_{\text{stuck}}(t) \rangle$ from the trajectories of each bead. Alternatively, if no stuck bead is present, and there are several (>5) beads in the FOV, an alternative approach consists in averaging the axial coordinates of all beads in the FOV to obtain the mean drift. Care must be taken to eliminate beads from this average that 504
505
506
507
508
509
510
511
512
513

514 display single-molecule events due to enzymatic activity. A
 515 second approach to eliminate drift has been recently developed
 516 that uses the Allan variance [4].

517 17. Single-molecule conditions are statistically achieved when the
 518 time between events (t_2) is considerably larger than the typical
 519 duration of a single event (t_1). The probability of two events
 520 occurring at the same time by chance is $(t_1/t_2)^2$. Thus, if
 521 $t_1/t_2=0.05$, the probability of two single-molecule events
 522 overlapping becomes $<1\%$.

523 **Acknowledgments**

524 We thank Francesco Pedaci and Antoine Le Gall for critical reading
 525 and very helpful comments. Financial support was provided by the
 526 Human Frontiers Science Program (M.N.) and the European
 527 Research Council (Starting Grant 260787 to M.N.).

528 **References**

- | | | |
|---|--|--|
| <p>529 1. te Velthuis AJW, Kerssemakers JWJ, Lipfert J
 530 et al (2010) Quantitative Guidelines for
 531 Force Calibration through Spectral Analysis
 532 of Magnetic Tweezers Data. <i>Biophys J</i> 99:
 533 1292–1302</p> <p>534 2. Lionnet T, Allemand JF, Revyakin A et al
 535 (2012) Single-molecule studies using magnetic
 536 traps. <i>Cold Spring Harb Protoc</i> 2012:34–49</p> <p>537 3. Vilfan ID, Lipfert J, Koster DA et al (2009)
 538 Magnetic Tweezers for Single-Molecule
 539 Experiments, <i>Handbook of Single-Molecule
 540 Biophysics</i>. Springer, New York</p> <p>541 4. Lansdorp BM, Saleh OA (2012) Power spec-
 542 trum and Allan variance methods for calibrat-
 543 ing single-molecule video-tracking instruments.
 544 <i>Rev Sci Instrum</i> 83:025115</p> <p>545 5. Strick T, Allemand J, Croquette V et al (2000)
 546 Twisting and stretching single DNA molecules.
 547 <i>Prog Biophys Mol Biol</i> 74:115–140</p> <p>548 6. Smith SB, Cui Y, Bustamante C (1996)
 549 Overstretching B-DNA: the elastic response of
 550 individual double-stranded and single-stranded
 551 DNA molecules. <i>Science</i> 271:795–799</p> <p>552 7. Bustamante C, Marko JF, Siggia ED et al
 553 (1994) Entropic elasticity of lambda-phage
 554 DNA. <i>Science</i> 265:1599–1600</p> <p>555 8. Strick TR, Allemand JF, Bensimon D et al
 556 (1996) The elasticity of a single supercoiled
 557 DNA molecule. <i>Science</i> 271:1835–1837</p> <p>558 9. Herrero-Galán E, Fuentes-Perez ME, Carrasco
 559 C et al (2013) Mechanical identities of RNA
 560 and DNA double helices unveiled at the single-
 561 molecule level. <i>J Am Chem Soc</i> 135:122–131</p> | <p>10. Abels JA, Moreno-Herrero F, van der Heijden
 T et al (2005) Single-molecule measurements
 of the persistence length of double-stranded
 RNA. <i>Biophys J</i> 88:2737–2744</p> <p>11. Liphardt J, Onoa B, Smith SB et al (2001)
 Reversible unfolding of single RNA molecules
 by mechanical force. <i>Science</i> 292:733–737</p> <p>12. Lipfert J, Wiggin M, Kerssemakers JWJ et al
 (2011) Freely orbiting magnetic tweezers to
 directly monitor changes in the twist of nucleic
 acids. <i>Nat Commun</i> 2:439</p> <p>13. Lebel P, Basu A, Oberstrass FC et al (2014)
 Gold rotor bead tracking for high-speed mea-
 surements of DNA twist, torque and extension.
 <i>Nat Methods</i> 11:456–462</p> <p>14. Bryant Z, Oberstrass FC, Basu A (2012) Recent
 developments in single-molecule DNA mechan-
 ics. <i>Curr Opin Struct Biol</i> 22:304–312</p> <p>15. Mosconi F, Allemand JF, Bensimon D et al
 (2009) Measurement of the torque on a single
 stretched and twisted DNA using magnetic
 tweezers. <i>Phys Rev Lett</i> 102:078301</p> <p>16. Crisona NJ, Strick TR, Bensimon D et al
 (2000) Preferential relaxation of positively
 supercoiled DNA by <i>E. coli</i> topoisomerase IV
 in single-molecule and ensemble measure-
 ments. <i>Genes Dev</i> 14:2881–2892</p> <p>17. Stone MD, Bryant Z, Crisona NJ et al (2003)
 Chirality sensing by <i>Escherichia coli</i> topoisomer-
 ase IV and the mechanism of type II topoisomer-
 ases. <i>Proc Natl Acad Sci U S A</i> 100:8654–8659</p> <p>18. Koster DA, Croquette V, Dekker C et al (2005)
 Friction and torque govern the relaxation of</p> | <p>562
563
564
565
566
567
568
569
570
571
572
573
574
575
576
577
578
579
580
581
582
583
584
585
586
587
588
589
590
591
592
593
594</p> |
|---|--|--|

- 595 DNA supercoils by eukaryotic topoisomerase 627
 596 IB. *Nature* 434:671–674 628
- 597 19. Nollmann M, Stone MD, Bryant Z et al (2007) 629
 598 Multiple modes of *Escherichia coli* DNA gyrase 630
 599 activity revealed by force and torque. *Nat*
 600 *Struct Mol Biol* 14:264–271 631
- 601 20. Ribbeck N, Saleh OA (2008) Multiplexed 632
 602 single-molecule measurements with magnetic 633
 603 tweezers. *Rev Sci Instrum* 79:094301 634
- 604 21. De Vlaminc I, Henighan T, van Loenhout 635
 605 MTJ et al (2011) Highly Parallel Magnetic 636
 606 Tweezers by Targeted DNA Tethering. *Nano*
 607 *Lett* 11:5489–5493 637
- 608 22. Ding F, Manosas M, Spiering MM et al (2012) 638
 609 Single-molecule mechanical identification and 639
 610 sequencing. *Nat Methods* 9:367–372 640
- 611 23. Manosas M, Perumal SK, Croquette V et al 641
 612 (2012) Direct observation of stalled fork 642
 613 restart via fork regression in the T4 replication 643
 614 system. *Science* 338:1217–1220 644
- 615 24. Manosas M, Meglio A, Spiering MM et al 645
 616 (2010) Magnetic Tweezers for the Study of 646
 617 DNA Tracking Motors. *Methods Enzymol* 647
 618 475:297–320 648
- 619 25. Lionnet T, Allemand JF, Revyakin A et al. 649
 620 (2011) Magnetic trap construction. *Cold*
 621 *Spring Harb Protoc* 2012, pdb.prot067496–
 622 pdb.prot067496 650
- 623 26. Revyakin A, Ebright RH, Strick TR (2005) 651
 624 Single-molecule DNA nanomanipulation: 652
 625 improved resolution through use of shorter 653
 626 DNA fragments. *Nat Methods* 2:127–138 654
27. Cheng W, Dumont S, Tinoco I et al (2007) 655
 NS3 helicase actively separates RNA strands 656
 and senses sequence barriers ahead of the open- 657
 ing fork. *Proc Natl Acad Sci U S A* 658
 104:13954–13959 659
28. Ovrzyn B, Izen S (2000) Imaging of transparent 660
 spheres through a planar interface using a high- 661
 numerical-aperture optical microscope. *J Opt*
 Soc Am A Opt Image Sci Vis 17:1202–1213 662
29. Gosse C, Croquette V (2002) Magnetic twee- 663
 zers: micromanipulation and force measurement 664
 at the molecular level. *Biophys J* 82:3314–3329 665
30. van Loenhout MTJ, Kerssemakers JWJ, De 666
 Vlaminc I et al (2012) Non-Bias-Limited 667
 Tracking of Spherical Particles, Enabling 668
 Nanometer Resolution at Low Magnification. 669
Biophys J 102:2362–2371 670
31. Lipfert J, Kerssemakers JJW, Rojer M et al 671
 (2011) A method to track rotational motion 672
 for use in single-molecule biophysics. *Rev Sci*
Instrum 82:103707 673
32. Cattoni DI, Fiche J-B, Valeri A et al (2013) 674
 Super-resolution imaging of bacteria in a 675
 microfluidics device. *PLoS One* 8:e76268 676
33. De Vlaminc I, Henighan T, van Loenhout 677
 MTJ et al (2012) Magnetic Forces and DNA 678
 Mechanics in Multiplexed Magnetic Tweezers. 679
PLoS One 7:e41432 680
34. Ptacin JL, Nollmann M, Becker EC et al (2008) 681
 Sequence-directed DNA export guides chromo- 682
 some translocation during sporulation in *Bacillus* 683
subtilis. *Nat Struct Mol Biol* 15:485–493 684